

Tensor modes in loop quantum cosmology

with G. Hossain, arXiv:0810.4330 [gr-qc]

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October 25th, 2008

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- To compute cosmological observables from LQC linear **tensor perturbations**.

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- To discuss related issues and future directions.

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$$\alpha = \frac{1+n}{3r} \lambda \left(\left| 1 + \frac{1}{\lambda} \right|^{\frac{3r}{2(1+n)}} - \left| 1 - \frac{1}{\lambda} \right|^{\frac{3r}{2(1+n)}} \right), \quad \lambda \sim \mathcal{V}^{2(1+n)/3}$$

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where

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Assuming $\Delta = \Delta_{\text{Pl}}$

$$4 < c \leq 6, \quad -0.01 \approx -\frac{1}{162} < \alpha_c < \frac{1}{9} \approx 0.1$$

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$$1 < q_\alpha < 6, \quad 1.6 \approx \frac{3^{3/4}}{\sqrt{2}} < \alpha_q < \frac{27}{4} \approx 6.8$$

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$$r = 1, \quad n = 1/2$$

$$c = 6, \quad \alpha_c = 0, \quad \alpha_q = \sqrt{3}, \quad q_\alpha = 3.$$

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and

$$\{\delta K_a^i(\mathbf{x}), \delta E_j^b(\mathbf{y})\} = 8\pi G \delta_a^b \delta_j^i \delta(\mathbf{x}, \mathbf{y})$$

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Only **inverse-volume corrections**:

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We solve it in large- and small-volume regimes separately.

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Inflation occurs for $p < -1$ (de Sitter: $p = -1$), **superinflation** when $-1 < p < 0$.

Outline

- 1 Background
- 2 Tensor perturbations
 - Near-Planckian regime
 - Quasi-classical regime

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- Large- and short-wavelength limits of the solution ($\nu > 0$)

$$w_k \sim -iC_1 \frac{2^\nu \Gamma(\nu)}{\pi} (-kz)^{1/2-\nu}, \quad |kz| \ll 1,$$

$$w_k \sim C_1 \sqrt{\frac{2}{\pi}} e^{-i(kz + \frac{\pi}{2}\nu + \frac{\pi}{4})}, \quad |kz| \gg 1.$$

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Plugging in the short-scale solution, one gets $|C_1| = \sqrt{8\pi^2 \ell_{\text{Pl}}^2/k}$.

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Tensor **spectrum**:

$$A_T^2 \equiv \frac{\mathcal{P}_h}{100} \equiv \frac{k^3}{200\pi^2 a^2} \sum_{+, \times} \langle |\hat{u}_{k \ll \mathcal{H}}|^2 \rangle \Big|_{k=k_*}$$

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Tensor spectral **index**:

$$n_T \equiv \left. \frac{d \ln A_T^2}{d \ln k} \right|_{k=k_*} = \frac{2(\epsilon + q_\alpha)}{\epsilon + q_\alpha - 1}$$

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- If $r \sim 10^{-8}$, still these bounds are $n_T < 1$.

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- **1**: r could be fine tuned to be small but scalar sector not available.
- **2**: Anomaly cancellation does not happen in scalar sector in this regime, which may be a sign that perturbation theory fails to converge.
- **3**: Close to the bounce, power-law evolution may not be a good approximation.

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- Near-Planckian phase might have occurred only at very early times (unobservably large scales) and for a short period.
- Scale-invariant or red-tilted tensor spectrum achieved in the interval $-1/q_\alpha < p \lesssim -1/(q_\alpha + 1)$, but could spoil scale invariance of **scalar spectrum**.
- **1**: r could be fine tuned to be small but scalar sector not available.
- **2**: Anomaly cancellation does not happen in scalar sector in this regime, which may be a sign that perturbation theory fails to converge.
- **3**: Close to the bounce, power-law evolution may not be a good approximation. However, $a \approx \text{const.}$

Outline

- 1 Background
- 2 Tensor perturbations
 - Near-Planckian regime
 - Quasi-classical regime

Quasi-classical regime: Solution

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Mukhanov equation:

$$\partial_{\tau}^2 w_k + c\mathcal{H}(\alpha - 1)\partial_{\tau} w_k + \{(2\alpha - 1)k^2 + \mathcal{H}^2[\epsilon - 2 - c(\alpha - 1)]\}w_k \approx 0.$$

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- Correction term decays in time.

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Tensor index:

$$n_T \approx 2(1 + p + cp\delta_{\text{Pl}}) = \frac{-2(\epsilon + c\delta_{\text{Pl}})}{1 - \epsilon}.$$

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- However, there are caveats to be addressed.
- Only nonperturbative formalisms (covariant, δN , separate universe, etc.) could be trusted (also relevant for **anomaly issue**).
- Quasi-classical result reliable, but scalar sector still under inspection.