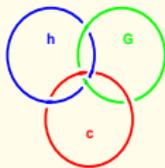


# Canonical Quantum General Relativity

Thomas Thiemann<sup>1,2,3</sup>

<sup>1</sup> Albert Einstein Institut, <sup>2</sup> FAU Erlangen – Nürnberg, <sup>3</sup> Perimeter Institute

PSU, Abhay Fest 2009



# Contents

- **Conceptual Foundations**
- Reduced Phase Space Quantisation
- Summary, Open Questions & Outlook

# Contents

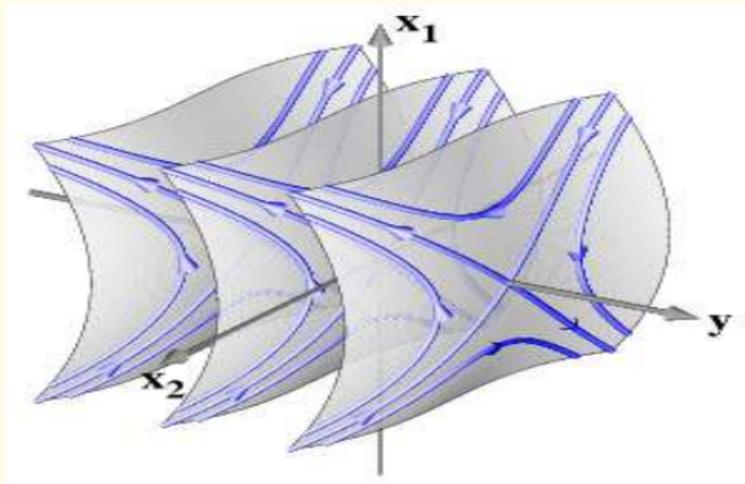
- Conceptual Foundations
- Reduced Phase Space Quantisation
- Summary, Open Questions & Outlook

# Contents

- Conceptual Foundations
- Reduced Phase Space Quantisation
- Summary, Open Questions & Outlook

# Classical Canonical Formulation

Canonical formulation:  $M \cong \mathbb{R} \times \sigma$



## Starting Point:

- Well posed (causal) initial value formulation for geometry and matter
- $\Rightarrow$  Globally hyperbolic spacetimes  $(M, g)$
- $\Rightarrow$  Topological restriction:  $M \cong \mathbb{R} \times \sigma$  [Geroch, 60's]
- No classical topology change, possibly quantum?

## Starting Point:

- Well posed (causal) initial value formulation for geometry and matter
- $\Rightarrow$  Globally hyperbolic spacetimes  $(M, g)$
- $\Rightarrow$  Topological restriction:  $M \cong \mathbb{R} \times \sigma$  [Geroch, 60's]
- No classical topology change, possibly quantum?

## Starting Point:

- Well posed (causal) initial value formulation for geometry and matter
- $\Rightarrow$  Globally hyperbolic spacetimes  $(M, g)$
- $\Rightarrow$  Topological restriction:  $M \cong \mathbb{R} \times \sigma$  [Geroch, 60's]
- No classical topology change, possibly quantum?

## Starting Point:

- Well posed (causal) initial value formulation for geometry and matter
- $\Rightarrow$  Globally hyperbolic spacetimes  $(M, g)$
- $\Rightarrow$  Topological restriction:  $M \cong \mathbb{R} \times \sigma$  [Geroch, 60's]
- No classical topology change, possibly quantum?

## Consequences: [ADM 60's]

- Consider arbitrary foliations  $Y : \mathbb{R} \times \sigma \rightarrow M$
- Require spacelike leaves of foliation  $\Sigma_t := Y(t, \sigma)$
- Pull all fields on  $M$  back to  $\mathbb{R} \times \sigma$
- Obtain velocity phase space of spatial fields (e.g. 3 – metric  $q_{ab}$  and extrinsic curvature  $K_{ab} \propto \partial q_{ab} / \partial t$ )
- Legendre transform  $K_{ab} \mapsto p^{ab}$  singular (due to  $\text{Diff}(M)$  invariance):  
Spatial diffeomorphism and Hamiltonian constraints  $c_a, c$
- Canonical Hamiltonian

$$H_{\text{canon}} = \int_{\sigma} d^3x \left( n c + v^a c_a \right) =: c(n) + \vec{c}(v)$$

## Consequences: [ADM 60's]

- Consider arbitrary foliations  $Y : \mathbb{R} \times \sigma \rightarrow M$
- Require spacelike leaves of foliation  $\Sigma_t := Y(t, \sigma)$
- Pull all fields on  $M$  back to  $\mathbb{R} \times \sigma$
- Obtain velocity phase space of spatial fields (e.g. 3 – metric  $q_{ab}$  and extrinsic curvature  $K_{ab} \propto \partial q_{ab} / \partial t$ )
- Legendre transform  $K_{ab} \mapsto p^{ab}$  singular (due to  $\text{Diff}(M)$  invariance):  
Spatial diffeomorphism and Hamiltonian constraints  $c_a, c$
- Canonical Hamiltonian

$$H_{\text{canon}} = \int_{\sigma} d^3x \left( n c + v^a c_a \right) =: c(n) + \vec{c}(v)$$

## Consequences: [ADM 60's]

- Consider arbitrary foliations  $Y : \mathbb{R} \times \sigma \rightarrow M$
- Require spacelike leaves of foliation  $\Sigma_t := Y(t, \sigma)$
- Pull all fields on  $M$  back to  $\mathbb{R} \times \sigma$
- Obtain velocity phase space of spatial fields (e.g. 3 – metric  $q_{ab}$  and extrinsic curvature  $K_{ab} \propto \partial q_{ab} / \partial t$ )
- Legendre transform  $K_{ab} \mapsto p^{ab}$  singular (due to  $\text{Diff}(M)$  invariance):  
Spatial diffeomorphism and Hamiltonian constraints  $c_a, c$
- Canonical Hamiltonian

$$H_{\text{canon}} = \int_{\sigma} d^3x \left( n c + v^a c_a \right) =: c(n) + \vec{c}(v)$$

## Consequences: [ADM 60's]

- Consider arbitrary foliations  $Y : \mathbb{R} \times \sigma \rightarrow M$
- Require spacelike leaves of foliation  $\Sigma_t := Y(t, \sigma)$
- Pull all fields on  $M$  back to  $\mathbb{R} \times \sigma$
- Obtain velocity phase space of spatial fields (e.g. 3 – metric  $q_{ab}$  and extrinsic curvature  $K_{ab} \propto \partial q_{ab} / \partial t$ )
- Legendre transform  $K_{ab} \mapsto p^{ab}$  singular (due to  $\text{Diff}(M)$  invariance):  
Spatial diffeomorphism and Hamiltonian constraints  $c_a, c$
- Canonical Hamiltonian

$$H_{\text{canon}} = \int_{\sigma} d^3x \left( n c + v^a c_a \right) =: c(n) + \vec{c}(v)$$

## Consequences: [ADM 60's]

- Consider arbitrary foliations  $Y : \mathbb{R} \times \sigma \rightarrow M$
- Require spacelike leaves of foliation  $\Sigma_t := Y(t, \sigma)$
- Pull all fields on  $M$  back to  $\mathbb{R} \times \sigma$
- Obtain velocity phase space of spatial fields (e.g. 3 – metric  $q_{ab}$  and extrinsic curvature  $K_{ab} \propto \partial q_{ab} / \partial t$ )
- Legendre transform  $K_{ab} \mapsto p^{ab}$  singular (due to  $\text{Diff}(M)$  invariance):  
Spatial diffeomorphism and Hamiltonian constraints  $c_a, c$
- Canonical Hamiltonian

$$H_{\text{canon}} = \int_{\sigma} d^3x \left( n c + v^a c_a \right) =: c(n) + \vec{c}(v)$$

## Consequences: [ADM 60's]

- Consider arbitrary foliations  $Y : \mathbb{R} \times \sigma \rightarrow M$
- Require spacelike leaves of foliation  $\Sigma_t := Y(t, \sigma)$
- Pull all fields on  $M$  back to  $\mathbb{R} \times \sigma$
- Obtain velocity phase space of spatial fields (e.g. 3 – metric  $q_{ab}$  and extrinsic curvature  $K_{ab} \propto \partial q_{ab} / \partial t$ )
- Legendre transform  $K_{ab} \mapsto p^{ab}$  singular (due to  $\text{Diff}(M)$  invariance):  
Spatial diffeomorphism and Hamiltonian constraints  $c_a, c$
- Canonical Hamiltonian

$$H_{\text{canon}} = \int_{\sigma} d^3x \left( n c + v^a c_a \right) =: c(n) + \vec{c}(v)$$

## Remarks:

- Algebraic structure of  $c$ ,  $c_a$  Foliation independent
- Symplectic structure of geometry and matter fields Foliation independent
- Foliation dependence encoded in lapse, shift  $n$ ,  $v^a$
- Foliation independence (Diff(M) invariance)  $\Rightarrow H_{\text{canon}} \approx 0$
- 10 Einstein Equations equivalent to

$$\partial_t q_{ab} = \{H_{\text{canon}}, q_{ab}\}, \quad \partial_t p^{ab} = \{H_{\text{canon}}, p^{ab}\}, \quad c = 0, \quad c_a = 0$$

- In particular, building  $g_{\mu\nu}$ ,  $n^\mu$  from  $q_{ab}$ ,  $n$ ,  $v^a$  one obtains  $q_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$  and

$$\{H_{\text{canon}}, q_{\mu\nu}\} = [\mathcal{L}_u q]_{\mu\nu}, \quad u^\mu = nn^\mu + (\vec{v})^\mu$$

recovery of Diff(M) (on shell).

## Remarks:

- Algebraic structure of  $c$ ,  $c_a$  Foliation independent
- Symplectic structure of geometry and matter fields Foliation independent
- Foliation dependence encoded in lapse, shift  $n$ ,  $v^a$
- Foliation independence (Diff(M) invariance)  $\Rightarrow H_{\text{canon}} \approx 0$
- 10 Einstein Equations equivalent to

$$\partial_t q_{ab} = \{H_{\text{canon}}, q_{ab}\}, \quad \partial_t p^{ab} = \{H_{\text{canon}}, p^{ab}\}, \quad c = 0, \quad c_a = 0$$

- In particular, building  $g_{\mu\nu}$ ,  $n^\mu$  from  $q_{ab}$ ,  $n$ ,  $v^a$  one obtains  $q_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$  and

$$\{H_{\text{canon}}, q_{\mu\nu}\} = [\mathcal{L}_u q]_{\mu\nu}, \quad u^\mu = nn^\mu + (\vec{v})^\mu$$

recovery of Diff(M) (on shell).

## Remarks:

- Algebraic structure of  $c$ ,  $c_a$  Foliation independent
- Symplectic structure of geometry and matter fields Foliation independent
- Foliation dependence encoded in lapse, shift  $n$ ,  $v^a$
- Foliation independence (Diff(M) invariance)  $\Rightarrow H_{\text{canon}} \approx 0$
- 10 Einstein Equations equivalent to

$$\partial_t q_{ab} = \{H_{\text{canon}}, q_{ab}\}, \quad \partial_t p^{ab} = \{H_{\text{canon}}, p^{ab}\}, \quad c = 0, \quad c_a = 0$$

- In particular, building  $g_{\mu\nu}$ ,  $n^\mu$  from  $q_{ab}$ ,  $n$ ,  $v^a$  one obtains  $q_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$  and

$$\{H_{\text{canon}}, q_{\mu\nu}\} = [\mathcal{L}_u q]_{\mu\nu}, \quad u^\mu = nn^\mu + (\vec{v})^\mu$$

recovery of Diff(M) (on shell).

## Remarks:

- Algebraic structure of  $c$ ,  $c_a$  Foliation independent
- Symplectic structure of geometry and matter fields Foliation independent
- Foliation dependence encoded in lapse, shift  $n$ ,  $v^a$
- Foliation independence (Diff(M) invariance)  $\Rightarrow H_{\text{canon}} \approx 0$
- 10 Einstein Equations equivalent to

$$\partial_t q_{ab} = \{H_{\text{canon}}, q_{ab}\}, \quad \partial_t p^{ab} = \{H_{\text{canon}}, p^{ab}\}, \quad c = 0, \quad c_a = 0$$

- In particular, building  $g_{\mu\nu}$ ,  $n^\mu$  from  $q_{ab}$ ,  $n$ ,  $v^a$  one obtains  $q_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$  and

$$\{H_{\text{canon}}, q_{\mu\nu}\} = [\mathcal{L}_u q]_{\mu\nu}, \quad u^\mu = nn^\mu + (\vec{v})^\mu$$

recovery of Diff(M) (on shell).

## Remarks:

- Algebraic structure of  $c$ ,  $c_a$  Foliation independent
- Symplectic structure of geometry and matter fields Foliation independent
- Foliation dependence encoded in lapse, shift  $n$ ,  $v^a$
- Foliation independence (Diff(M) invariance)  $\Rightarrow H_{\text{canon}} \approx 0$
- 10 Einstein Equations equivalent to

$$\partial_t q_{ab} = \{H_{\text{canon}}, q_{ab}\}, \quad \partial_t p^{ab} = \{H_{\text{canon}}, p^{ab}\}, \quad c = 0, \quad c_a = 0$$

- In particular, building  $g_{\mu\nu}$ ,  $n^\mu$  from  $q_{ab}$ ,  $n$ ,  $v^a$  one obtains  $q_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$  and

$$\{H_{\text{canon}}, q_{\mu\nu}\} = [\mathcal{L}_u q]_{\mu\nu}, \quad u^\mu = nn^\mu + (\vec{v})^\mu$$

recovery of Diff(M) (on shell).

## Remarks:

- Algebraic structure of  $c, c_a$  Foliation independent
- Symplectic structure of geometry and matter fields Foliation independent
- Foliation dependence encoded in lapse, shift  $n, v^a$
- Foliation independence (Diff(M) invariance)  $\Rightarrow H_{\text{canon}} \approx 0$
- 10 Einstein Equations equivalent to

$$\partial_t q_{ab} = \{H_{\text{canon}}, q_{ab}\}, \quad \partial_t p^{ab} = \{H_{\text{canon}}, p^{ab}\}, \quad c = 0, \quad c_a = 0$$

- In particular, building  $g_{\mu\nu}, n^\mu$  from  $q_{ab}, n, v^a$  one obtains  $q_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu$  and

$$\{H_{\text{canon}}, q_{\mu\nu}\} = [\mathcal{L}_u q]_{\mu\nu}, \quad u^\mu = nn^\mu + (\vec{v})^\mu$$

recovery of Diff(M) (on shell).

## Consistency:

- First class (Dirac) hypersurface deformation algebra  $\mathfrak{D}$

$$\begin{aligned}\{\vec{c}(v), \vec{c}(v')\} &= -\vec{c}([v, v']) \\ \{\vec{c}(v), c(n)\} &= -c(v[n]) \\ \{c(n), c(n')\} &= -\vec{c}(q^{-1}[n \, dn' - n' \, dn])\end{aligned}$$

- **Universality:** purely geometric origin, independent of matter content  
[Hojman, Kuchař, Teitelboim 70's]
- spatial diffeos generate subalgebra but not ideal
- $\mathfrak{D}$  no Lie algebra (structure functions)

## Consistency:

- First class (Dirac) hypersurface deformation algebra  $\mathfrak{D}$

$$\begin{aligned}\{\vec{c}(v), \vec{c}(v')\} &= -\vec{c}([v, v']) \\ \{\vec{c}(v), c(n)\} &= -c(v[n]) \\ \{c(n), c(n')\} &= -\vec{c}(q^{-1}[n \, dn' - n' \, dn])\end{aligned}$$

- **Universality:** purely geometric origin, independent of matter content  
[Hojman, Kuchař, Teitelboim 70's]
- spatial diffeos generate subalgebra but not ideal
- $\mathfrak{D}$  no Lie algebra (structure functions)

## Consistency:

- First class (Dirac) hypersurface deformation algebra  $\mathfrak{D}$

$$\begin{aligned}\{\vec{c}(v), \vec{c}(v')\} &= -\vec{c}([v, v']) \\ \{\vec{c}(v), c(n)\} &= -c(v[n]) \\ \{c(n), c(n')\} &= -\vec{c}(q^{-1}[n \, dn' - n' \, dn])\end{aligned}$$

- **Universality:** purely geometric origin, independent of matter content  
[Hojman, Kuchař, Teitelboim 70's]
- spatial diffeos generate subalgebra but not ideal
- $\mathfrak{D}$  no Lie algebra (structure functions)

## Consistency:

- First class (Dirac) hypersurface deformation algebra  $\mathfrak{D}$

$$\begin{aligned}\{\vec{c}(v), \vec{c}(v')\} &= -\vec{c}([v, v']) \\ \{\vec{c}(v), c(n)\} &= -c(v[n]) \\ \{c(n), c(n')\} &= -\vec{c}(q^{-1}[n \, dn' - n' \, dn])\end{aligned}$$

- **Universality:** purely geometric origin, independent of matter content  
[Hojman, Kuchař, Teitelboim 70's]
- spatial diffeos generate subalgebra but not ideal
- $\mathfrak{D}$  no Lie algebra (structure functions)

# Problem of Time

## Interpretation:

- $H_{\text{canon}}$  constrained to vanish, no true Hamiltonian
- $H_{\text{canon}}$  generates gauge transformations, not physical evolution
- $q_{ab}, p^{ab}, \dots$  not gauge invariant, not observable
- $\{H_{\text{canon}}, O\} = 0$  for observable, gauge invariant  $O$
- **Problem of time:** Dynamical interpretation?

# Problem of Time

## Interpretation:

- $H_{\text{canon}}$  constrained to vanish, no true Hamiltonian
- $H_{\text{canon}}$  generates gauge transformations, not physical evolution
- $q_{ab}, p^{ab}, \dots$  not gauge invariant, not observable
- $\{H_{\text{canon}}, O\} = 0$  for observable, gauge invariant  $O$
- **Problem of time:** Dynamical interpretation?

# Problem of Time

## Interpretation:

- $H_{\text{canon}}$  constrained to vanish, no true Hamiltonian
- $H_{\text{canon}}$  generates gauge transformations, not physical evolution
- $q_{ab}, p^{ab}, \dots$  not gauge invariant, not observable
- $\{H_{\text{canon}}, O\} = 0$  for observable, gauge invariant  $O$
- **Problem of time:** Dynamical interpretation?

# Problem of Time

## Interpretation:

- $H_{\text{canon}}$  constrained to vanish, no true Hamiltonian
- $H_{\text{canon}}$  generates gauge transformations, not physical evolution
- $q_{ab}, p^{ab}, \dots$  not gauge invariant, not observable
- $\{H_{\text{canon}}, O\} = 0$  for observable, gauge invariant  $O$
- **Problem of time:** Dynamical interpretation?

# Problem of Time

## Interpretation:

- $H_{\text{canon}}$  constrained to vanish, no true Hamiltonian
- $H_{\text{canon}}$  generates gauge transformations, not physical evolution
- $q_{ab}, p^{ab}, \dots$  not gauge invariant, not observable
- $\{H_{\text{canon}}, O\} = 0$  for observable, gauge invariant  $O$
- **Problem of time:** Dynamical interpretation?

## Gravitational Higgs Mechanism [Brown & Kuchař 90's]

- In GR, gauge invariant definition of curvature etc. only relative to geodesic test observers [Wald 90's]
- Test observers = mathematical idealisation
- Brown – Kuchař dust action: 4 scalar fields  $T, S^J$  minimally coupled however: geometry backreaction taken seriously
- Natural: Superposition of  $\infty$  # of point particle actions
- EL Equations: Dust particles move on unit geodesics,  $T(x)$  = proper time along geodesic through  $x$ ,  $S^J(x)$  labels geodesic
- Cold Dark Matter candidate (NIMP)

## Gravitational Higgs Mechanism [Brown & Kuchař 90's]

- In GR, gauge invariant definition of curvature etc. only relative to geodesic test observers [Wald 90's]
- Test observers = mathematical idealisation
- Brown – Kuchař dust action: 4 scalar fields  $T, S^J$  minimally coupled however: geometry backreaction taken seriously
- Natural: Superposition of  $\infty$  # of point particle actions
- EL Equations: Dust particles move on unit geodesics,  $T(x)$  = proper time along geodesic through  $x$ ,  $S^J(x)$  labels geodesic
- Cold Dark Matter candidate (NIMP)

## Gravitational Higgs Mechanism [Brown & Kuchař 90's]

- In GR, gauge invariant definition of curvature etc. only relative to geodesic test observers [Wald 90's]
- Test observers = mathematical idealisation
- Brown – Kuchař dust action: 4 scalar fields  $T, S^J$  minimally coupled however: geometry backreaction taken seriously
- Natural: Superposition of  $\infty$  # of point particle actions
- EL Equations: Dust particles move on unit geodesics,  $T(x)$  = proper time along geodesic through  $x$ ,  $S^J(x)$  labels geodesic
- Cold Dark Matter candidate (NIMP)

## Gravitational Higgs Mechanism [Brown & Kuchař 90's]

- In GR, gauge invariant definition of curvature etc. only relative to geodesic test observers [Wald 90's]
- Test observers = mathematical idealisation
- Brown – Kuchař dust action: 4 scalar fields  $T, S^J$  minimally coupled however: geometry backreaction taken seriously
- Natural: Superposition of  $\infty$  # of point particle actions
- EL Equations: Dust particles move on unit geodesics,  $T(x)$  = proper time along geodesic through  $x$ ,  $S^J(x)$  labels geodesic
- Cold Dark Matter candidate (NIMP)

## Gravitational Higgs Mechanism [Brown & Kuchař 90's]

- In GR, gauge invariant definition of curvature etc. only relative to geodesic test observers [Wald 90's]
- Test observers = mathematical idealisation
- Brown – Kuchař dust action: 4 scalar fields  $T, S^J$  minimally coupled however: geometry backreaction taken seriously
- Natural: Superposition of  $\infty$  # of point particle actions
- EL Equations: Dust particles move on unit geodesics,  $T(x)$  = proper time along geodesic through  $x$ ,  $S^J(x)$  labels geodesic
- Cold Dark Matter candidate (NIMP)

## Gravitational Higgs Mechanism [Brown & Kuchař 90's]

- In GR, gauge invariant definition of curvature etc. only relative to geodesic test observers [Wald 90's]
- Test observers = mathematical idealisation
- Brown – Kuchař dust action: 4 scalar fields  $T, S^J$  minimally coupled however: geometry backreaction taken seriously
- Natural: Superposition of  $\infty$  # of point particle actions
- EL Equations: Dust particles move on unit geodesics,  $T(x)$  = proper time along geodesic through  $x$ ,  $S^J(x)$  labels geodesic
- Cold Dark Matter candidate (NIMP)

● Deparametrisation:

$$\mathbf{c} := \mathbf{c}^D + \mathbf{c}^{ND}, \quad c_a = c_a^D + c_a^{ND} \Rightarrow \tilde{\mathbf{c}} = \mathbf{P} + \mathbf{h}, \quad \mathbf{h} = \sqrt{[\mathbf{c}^{ND}]^2 - q^{ab} c_a^{ND} c_b^{ND}}$$

● For close to flat geometry  $\mathbf{h} \approx \mathbf{c}^{ND} \approx \mathbf{h}^{SM}$  **hard to achieve!**

● Remarkably  $\{\tilde{\mathbf{c}}(n), \tilde{\mathbf{c}}(n')\} = 0$  [Brown & Kuchař 90's]

⇒ **Explicit** relational solution [Bergmann 60's, Rovelli 90's, Dittrich 00's]

● First symplectic reduction wrt  $c_a$  [Kuchař 90's] e.g.

$$q_{ab}(x) \rightarrow q_{JK}(s) := [q_{ab}(x) S_J^a(x) S_K^b(x)]_{S^J(x)=s^J}, \quad S_J^a S_{,b}^J = \delta_b^a, \quad S_J^a S_{,a}^K = \delta_J^K$$

● For any spatially diffeo inv., dust indep.  $f$  get observable

$$O_f(\tau) := \exp(\{H_\tau, \cdot\}) \cdot f, \quad H_\tau := \int_\sigma d^3x (\tau - T(x)) h^{ND}(x)$$

● Physical time evolution

$$\frac{d}{d\tau} O_f(\tau) = \{H_{\text{phys}}, O_f(\tau)\}, \quad H_{\text{phys}} := \int_\sigma d^3x h^{ND}(x)$$

- Deparametrisation:

$$c := c^D + c^{ND}, \quad c_a = c_a^D + c_a^{ND} \Rightarrow \tilde{c} = P + h, \quad h = \sqrt{[c^{ND}]^2 - q^{ab}c_a^{ND}c_b^{ND}}$$

- For close to flat geometry  $h \approx c^{ND} \approx h^{SM}$  **hard to achieve!**

- Remarkably  $\{\tilde{c}(n), \tilde{c}(n')\} = 0$  [Brown & Kuchař 90's]

$\Rightarrow$  **Explicit** relational solution [Bergmann 60's, Rovelli 90's, Dittrich 00's]

- First symplectic reduction wrt  $c_a$  [Kuchař 90's] e.g.

$$q_{ab}(x) \rightarrow q_{JK}(s) := [q_{ab}(x)S_J^a(x)S_K^b(x)]_{S^J(x)=s^J}, \quad S_J^a S_{,b}^J = \delta_b^a, \quad S_J^a S_{,a}^K = \delta_J^K$$

- For any spatially diffeo inv., dust indep.  $f$  get observable

$$O_f(\tau) := \exp(\{H_\tau, \cdot\}) \cdot f, \quad H_\tau := \int_\sigma d^3x (\tau - T(x)) h^{ND}(x)$$

- Physical time evolution

$$\frac{d}{d\tau} O_f(\tau) = \{H_{\text{phys}}, O_f(\tau)\}, \quad H_{\text{phys}} := \int_\sigma d^3x h^{ND}(x)$$

- Deparametrisation:

$$c := c^D + c^{ND}, \quad c_a = c_a^D + c_a^{ND} \Rightarrow \tilde{c} = P + h, \quad h = \sqrt{[c^{ND}]^2 - q^{ab}c_a^{ND}c_b^{ND}}$$

- For close to flat geometry  $h \approx c^{ND} \approx h^{SM}$  **hard to achieve!**

- Remarkably  $\{\tilde{c}(n), \tilde{c}(n')\} = 0$  [Brown & Kuchař 90's]

⇒ **Explicit** relational solution [Bergmann 60's, Rovelli 90's, Dittrich 00's]

- First symplectic reduction wrt  $c_a$  [Kuchař 90's] e.g.

$$q_{ab}(x) \rightarrow q_{JK}(s) := [q_{ab}(x)S_J^a(x)S_K^b(x)]_{S^J(x)=s^J}, \quad S_J^a S_{,b}^J = \delta_b^a, \quad S_J^a S_{,a}^K = \delta_J^K$$

- For any spatially diffeo inv., dust indep.  $f$  get observable

$$O_f(\tau) := \exp(\{H_\tau, \cdot\}) \cdot f, \quad H_\tau := \int_\sigma d^3x (\tau - T(x)) h^{ND}(x)$$

- Physical time evolution

$$\frac{d}{d\tau} O_f(\tau) = \{H_{\text{phys}}, O_f(\tau)\}, \quad H_{\text{phys}} := \int_\sigma d^3x h^{ND}(x)$$

- Deparametrisation:

$$c := c^D + c^{ND}, \quad c_a = c_a^D + c_a^{ND} \Rightarrow \tilde{c} = P + h, \quad h = \sqrt{[c^{ND}]^2 - q^{ab}c_a^{ND}c_b^{ND}}$$

- For close to flat geometry  $h \approx c^{ND} \approx h^{SM}$  **hard to achieve!**
- Remarkably  $\{\tilde{c}(n), \tilde{c}(n')\} = 0$  [Brown & Kuchař 90's]  
 $\Rightarrow$  **Explicit** relational solution [Bergmann 60's, Rovelli 90's, Dittrich 00's]
- First symplectic reduction wrt  $c_a$  [Kuchař 90's] e.g.

$$q_{ab}(x) \rightarrow q_{JK}(s) := [q_{ab}(x)S_J^a(x)S_K^b(x)]_{S^J(x)=s^J}, \quad S_J^a S_{,b}^J = \delta_b^a, \quad S_J^a S_{,a}^K = \delta_J^K$$

- For any spatially diffeo inv., dust indep.  $f$  get observable

$$O_f(\tau) := \exp(\{H_\tau, \cdot\}) \cdot f, \quad H_\tau := \int_\sigma d^3x (\tau - T(x)) h^{ND}(x)$$

- Physical time evolution

$$\frac{d}{d\tau} O_f(\tau) = \{H_{\text{phys}}, O_f(\tau)\}, \quad H_{\text{phys}} := \int_\sigma d^3x h^{ND}(x)$$

- Deparametrisation:

$$c := c^D + c^{ND}, \quad c_a = c_a^D + c_a^{ND} \Rightarrow \tilde{c} = P + h, \quad h = \sqrt{[c^{ND}]^2 - q^{ab}c_a^{ND}c_b^{ND}}$$

- For close to flat geometry  $h \approx c^{ND} \approx h^{SM}$  **hard to achieve!**
- Remarkably  $\{\tilde{c}(n), \tilde{c}(n')\} = 0$  [Brown & Kuchař 90's]  
 $\Rightarrow$  **Explicit** relational solution [Bergmann 60's, Rovelli 90's, Dittrich 00's]
- First symplectic reduction wrt  $c_a$  [Kuchař 90's] e.g.

$$q_{ab}(x) \rightarrow q_{JK}(s) := [q_{ab}(x)S_J^a(x)S_K^b(x)]_{S^J(x)=s^J}, \quad S_J^a S_{,b}^J = \delta_b^a, \quad S_J^a S_{,a}^K = \delta_J^K$$

- For any spatially diffeo inv., dust indep.  $f$  get observable

$$O_f(\tau) := \exp(\{H_\tau, \cdot\}) \cdot f, \quad H_\tau := \int_\sigma d^3x (\tau - T(x)) h^{ND}(x)$$

- Physical time evolution

$$\frac{d}{d\tau} O_f(\tau) = \{H_{\text{phys}}, O_f(\tau)\}, \quad H_{\text{phys}} := \int_\sigma d^3x h^{ND}(x)$$

- Deparametrisation:

$$c := c^D + c^{ND}, \quad c_a = c_a^D + c_a^{ND} \Rightarrow \tilde{c} = P + h, \quad h = \sqrt{[c^{ND}]^2 - q^{ab}c_a^{ND}c_b^{ND}}$$

- For close to flat geometry  $h \approx c^{ND} \approx h^{SM}$  **hard to achieve!**
- Remarkably  $\{\tilde{c}(n), \tilde{c}(n')\} = 0$  [Brown & Kuchař 90's]  
 $\Rightarrow$  **Explicit** relational solution [Bergmann 60's, Rovelli 90's, Dittrich 00's]
- First symplectic reduction wrt  $c_a$  [Kuchař 90's] e.g.

$$q_{ab}(x) \rightarrow q_{JK}(s) := [q_{ab}(x)S_J^a(x)S_K^b(x)]_{S^J(x)=s^J}, \quad S_J^a S_{,b}^J = \delta_b^a, \quad S_J^a S_{,a}^K = \delta_J^K$$

- For any spatially diffeo inv., dust indep.  $f$  get observable

$$O_f(\tau) := \exp(\{H_\tau, \cdot\}) \cdot f, \quad H_\tau := \int_\sigma d^3x (\tau - T(x)) h^{ND}(x)$$

- Physical time evolution

$$\frac{d}{d\tau} O_f(\tau) = \{H_{\text{phys}}, O_f(\tau)\}, \quad H_{\text{phys}} := \int_\sigma d^3x h^{ND}(x)$$

- Closed observable algebra due to automorphism property of Hamiltonian flow

$$\{O_f(\tau), O_{f'}(\tau)\} = O_{\{f, f'\}}(\tau)$$

- Reduced phase space Q'ion conceivable since e.g.

$$Q_{JK}(s) := O_{q_{JK}(s)}(0), \quad P^{JK}(s) := O_{p^{JK}(s)}(0)$$

$$\Rightarrow \{P^{JK}(s), Q_{LM}(s')\} = \delta_L^{(J} \delta_M^{K)} \delta(s, s')$$

- Closed observable algebra due to automorphism property of Hamiltonian flow

$$\{O_f(\tau), O_{f'}(\tau)\} = O_{\{f, f'\}}(\tau)$$

- Reduced phase space Q'ion conceivable since e.g.

$$Q_{JK}(s) := O_{q_{JK}(s)}(0), \quad P^{JK}(s) := O_{p^{JK}(s)}(0)$$

$$\Rightarrow \{P^{JK}(s), Q_{LM}(s')\} = \delta_L^{(J} \delta_M^{K)} \delta(s, s')$$

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c_j^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c_j^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c_j^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

## Physics of the Dust:

- Dust = Gravitational **Higgs**, Non-Dust = Gravitational **Goldstone Bosons**
- Conservative Hamiltonian system w/o constraints but true Hamiltonian
- Hamiltonian EOM wrt  $H_{\text{phys}}$  of physical Non-Dust dof agree with Gauge Transformations wrt  $H_{\text{canon}}$  of unphysical Non-Dust dof under proper field substitutions, e.g.  $q_{ab}(x) \leftrightarrow Q^{jk}(s)$
- No constraints but energy – momentum current conservation law

$$\{H_{\text{phys}}, O_{h^{\text{ND}}(s)}\} = 0, \quad \{H_{\text{phys}}, O_{c_j^{\text{ND}}(s)}\} = 0,$$

- Effectively **decouples** 4 Goldstone modes, agreement with observation (gravitational waves) [Giesel, Hofmann, T.T., Winkler 00's], [Giesel, Tambornino, T.T. 00's]
- In terms of  $\tilde{c}$  dust fields are **perfect** (nowhere singular) clocks
- Effective Action displays similarities with **Hořava action**

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\hat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reps. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reps. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\hat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reprs. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reprs. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\hat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reprs. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reprs. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\hat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reps. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reps. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\hat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reps. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reps. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\widehat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reps. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reps. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\widehat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reps. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reps. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathfrak{D}$

# Canonical Quantisation Strategies

- **Objective:** Irreducible representation of the  $*$ -algebra (or  $C^*$ )  $\mathfrak{A}_{\text{phys}}$  of Dirac observables supporting  $\widehat{H}_{\text{phys}}$
- Strategy 1: Constraint Q'ion (CQ) = Q'ion before reduction
- Strategy 2: Reduced phase space Q'ion (RQ) = Q'ion after reduction
- Complementary Advantages and Disadvantages
  - CQ+: Reps. of  $\mathfrak{A}_{\text{kin}}$  easy to find
  - CQ-: Phys. HS = Kernel(constraints) construction complicated (group averaging)
  - RQ+: Directly phys. HS w/o redundant dof in  $\mathfrak{A}_{\text{kin}}$
  - RQ-: Reps. of  $\mathfrak{A}_{\text{phys}}$  often difficult to find
- With dust, reduced phase space q'ion simpler, avoid difficult representation of  $\mathcal{D}$

# Algebra of Kinematical Functions

## Gauge Theory Formulation:

- Due to fermionic dof need to start with Palatini/Holst action [Ashtekar 80's], [Barbero, Holst, Immirzi 90's]
- After solving 2nd class (simplicity) constraints obtain

$$\{E_j^a(x), A_b^k(y)\} = \kappa \delta_b^a \delta_j^k \delta(x, y)$$

- Non-dust, gravitational contributions to the constraints

$$\begin{aligned} C_j^{\text{geo}} &= \mathcal{D}_a E_j^a \\ C_a^{\text{geo}} &= \text{Tr}(F_{ab} E^b) \\ C^{\text{geo}} &= \frac{\text{Tr}(F_{ab}[E^a, E^b])}{\sqrt{|\det(E)|}} + \dots \end{aligned}$$

# Algebra of Kinematical Functions

## Gauge Theory Formulation:

- Due to fermionic dof need to start with Palatini/Holst action [Ashtekar 80's], [Barbero, Holst, Immirzi 90's]
- After solving 2nd class (simplicity) constraints obtain

$$\{E_j^a(\mathbf{x}), A_b^k(\mathbf{y})\} = \kappa \delta_b^a \delta_j^k \delta(\mathbf{x}, \mathbf{y})$$

- Non-dust, gravitational contributions to the constraints

$$\begin{aligned} C_j^{\text{geo}} &= \mathcal{D}_a E_j^a \\ C_a^{\text{geo}} &= \text{Tr}(F_{ab} E^b) \\ C^{\text{geo}} &= \frac{\text{Tr}(F_{ab}[E^a, E^b])}{\sqrt{|\det(E)|}} + \dots \end{aligned}$$

# Algebra of Kinematical Functions

## Gauge Theory Formulation:

- Due to fermionic dof need to start with Palatini/Holst action [Ashtekar 80's], [Barbero, Holst, Immirzi 90's]
- After solving 2nd class (simplicity) constraints obtain

$$\{E_j^a(\mathbf{x}), A_b^k(\mathbf{y})\} = \kappa \delta_b^a \delta_j^k \delta(\mathbf{x}, \mathbf{y})$$

- Non-dust, gravitational contributions to the constraints

$$\begin{aligned} C_j^{\text{geo}} &= \mathcal{D}_a E_j^a \\ C_a^{\text{geo}} &= \text{Tr}(F_{ab} E^b) \\ C^{\text{geo}} &= \frac{\text{Tr}(F_{ab}[E^a, E^b])}{\sqrt{|\det(E)|}} + \dots \end{aligned}$$

# Algebra of Physical Observables

- Simply define (similar for  $E_i^j(\mathbf{s})$ )

$$A_i^j(\mathbf{s}) := O_{a_i^j(\mathbf{s})}(0), \quad a_i^j(\mathbf{s}) := [A_a^j S_i^a](x)_{S(x)=\mathbf{s}},$$

- Then

$$\{E_i^j(\mathbf{s}), A_j^k(\mathbf{s}')\} = \kappa \delta_j^k \delta_i^l \delta(\mathbf{s}, \mathbf{s}')$$

- No constraints but phys. Hamiltonian ( $\Sigma = S(\sigma)$ )

$$H = \int_{\Sigma} \sqrt{|-\eta^{\mu\nu} \text{Tr}(\tau_{\mu} F \wedge \{A, V\}) \text{Tr}(\tau_{\nu} F \wedge \{A, V\})|} =: \int d^3s H(\mathbf{s})$$

- Physical total volume

$$V = \int_{\Sigma} \sqrt{|\det(E)|}$$

- Symmetry group of H:  $\mathcal{G} = \mathcal{N} \times \text{Diff}(\Sigma)$
- $\mathcal{N}$ : Abelian normal subgroup generated by H(s), **active**  $\text{Diff}(\Sigma)$

# Algebra of Physical Observables

- Simply define (similar for  $E_i^j(\mathbf{s})$ )

$$A_i^j(\mathbf{s}) := O_{a_i^j(\mathbf{s})}(0), \quad a_i^j(\mathbf{s}) := [A_a^j S_i^a](x)_{S(x)=\mathbf{s}},$$

- Then

$$\{E_i^j(\mathbf{s}), A_j^k(\mathbf{s}')\} = \kappa \delta_i^k \delta_j^l \delta(\mathbf{s}, \mathbf{s}')$$

- No constraints but phys. Hamiltonian ( $\Sigma = S(\sigma)$ )

$$H = \int_{\Sigma} \sqrt{|-\eta^{\mu\nu} \text{Tr}(\tau_{\mu} F \wedge \{A, V\}) \text{Tr}(\tau_{\nu} F \wedge \{A, V\})|} =: \int d^3s H(\mathbf{s})$$

- Physical total volume

$$V = \int_{\Sigma} \sqrt{|\det(E)|}$$

- Symmetry group of H:  $\mathcal{G} = \mathcal{N} \times \text{Diff}(\Sigma)$

- $\mathcal{N}$ : Abelian normal subgroup generated by H(s), **active**  $\text{Diff}(\Sigma)$

# Algebra of Physical Observables

- Simply define (similar for  $E_i^j(\mathbf{s})$ )

$$A_i^j(\mathbf{s}) := O_{a_i^j(\mathbf{s})}(0), \quad a_i^j(\mathbf{s}) := [A_a^j S_i^a](x)_{S(x)=\mathbf{s}},$$

- Then

$$\{E_i^j(\mathbf{s}), A_j^k(\mathbf{s}')\} = \kappa \delta_i^k \delta_j^l \delta(\mathbf{s}, \mathbf{s}')$$

- No constraints but phys. Hamiltonian ( $\Sigma = S(\sigma)$ )

$$H = \int_{\Sigma} \sqrt{|-\eta^{\mu\nu} \text{Tr}(\tau_{\mu} \mathbf{F} \wedge \{\mathbf{A}, \mathbf{V}\}) \text{Tr}(\tau_{\nu} \mathbf{F} \wedge \{\mathbf{A}, \mathbf{V}\})|} =: \int d^3\mathbf{s} H(\mathbf{s})$$

- Physical total volume

$$V = \int_{\Sigma} \sqrt{|\det(\mathbf{E})|}$$

- Symmetry group of H:  $\mathcal{G} = \mathcal{N} \times \text{Diff}(\Sigma)$
- $\mathcal{N}$ : Abelian normal subgroup generated by H(s), **active**  $\text{Diff}(\Sigma)$

# Algebra of Physical Observables

- Simply define (similar for  $E_i^j(\mathbf{s})$ )

$$A_i^j(\mathbf{s}) := O_{a_i^j(\mathbf{s})}(0), \quad a_i^j(\mathbf{s}) := [A_a^j S_i^a](x)_{S(x)=\mathbf{s}},$$

- Then

$$\{E_i^j(\mathbf{s}), A_j^k(\mathbf{s}')\} = \kappa \delta_j^k \delta_i^l \delta(\mathbf{s}, \mathbf{s}')$$

- No constraints but phys. Hamiltonian ( $\Sigma = S(\sigma)$ )

$$H = \int_{\Sigma} \sqrt{|-\eta^{\mu\nu} \text{Tr}(\tau_{\mu} F \wedge \{A, V\}) \text{Tr}(\tau_{\nu} F \wedge \{A, V\})|} =: \int d^3s H(\mathbf{s})$$

- Physical total volume

$$V = \int_{\Sigma} \sqrt{|\det(E)|}$$

- Symmetry group of H:  $\mathcal{G} = \mathcal{N} \times \text{Diff}(\Sigma)$
- $\mathcal{N}$ : Abelian normal subgroup generated by H(s), **active**  $\text{Diff}(\Sigma)$

# Algebra of Physical Observables

- Simply define (similar for  $E_i^j(\mathbf{s})$ )

$$A_i^j(\mathbf{s}) := O_{a_i^j(\mathbf{s})}(0), \quad a_i^j(\mathbf{s}) := [A_a^j S_i^a](x)_{S(x)=\mathbf{s}},$$

- Then

$$\{E_i^j(\mathbf{s}), A_j^k(\mathbf{s}')\} = \kappa \delta_j^k \delta_i^l \delta(\mathbf{s}, \mathbf{s}')$$

- No constraints but phys. Hamiltonian ( $\Sigma = S(\sigma)$ )

$$H = \int_{\Sigma} \sqrt{|-\eta^{\mu\nu} \text{Tr}(\tau_{\mu} F \wedge \{A, V\}) \text{Tr}(\tau_{\nu} F \wedge \{A, V\})|} =: \int d^3s H(\mathbf{s})$$

- Physical total volume

$$V = \int_{\Sigma} \sqrt{|\det(\mathbf{E})|}$$

- Symmetry group of H:  $\mathfrak{G} = \mathcal{N} \times \text{Diff}(\Sigma)$

- $\mathcal{N}$ : Abelian normal subgroup generated by H(s), active  $\text{Diff}(\Sigma)$

# Algebra of Physical Observables

- Simply define (similar for  $E_i^j(\mathbf{s})$ )

$$A_i^j(\mathbf{s}) := O_{a_i^j(\mathbf{s})}(0), \quad a_i^j(\mathbf{s}) := [A_a^j S_i^a](x)_{S(x)=\mathbf{s}},$$

- Then

$$\{E_i^j(\mathbf{s}), A_j^k(\mathbf{s}')\} = \kappa \delta_i^k \delta_j^l \delta(\mathbf{s}, \mathbf{s}')$$

- No constraints but phys. Hamiltonian ( $\Sigma = S(\sigma)$ )

$$H = \int_{\Sigma} \sqrt{|-\eta^{\mu\nu} \text{Tr}(\tau_{\mu} \mathbf{F} \wedge \{\mathbf{A}, \mathbf{V}\}) \text{Tr}(\tau_{\nu} \mathbf{F} \wedge \{\mathbf{A}, \mathbf{V}\})|} =: \int d^3\mathbf{s} H(\mathbf{s})$$

- Physical total volume

$$V = \int_{\Sigma} \sqrt{|\det(\mathbf{E})|}$$

- Symmetry group of H:  $\mathfrak{G} = \mathcal{N} \times \text{Diff}(\Sigma)$
- $\mathcal{N}$ : Abelian normal subgroup generated by H(s), **active** Diff( $\Sigma$ )

# Physical Hilbert Space

**Lattice – inspired canon. gauge theory variables** [Gambini & Trias 81], [Jacobson, Rovelli, Smolin 88]

- Magnet. dof.: Holonomy (Wilson – Loop)

$$A(e) = \mathcal{P} \exp\left(\int_e A\right)$$

- Electr. dof: flux

$$E_j(S) = \int_S \epsilon_{abc} E_j^a dx^b \wedge dx^c$$

- Poisson – brackets:

$$\{E_j(S), A(e)\} = G A(e_1) \tau_j A(e_2); \quad e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$

# Physical Hilbert Space

**Lattice – inspired canon. gauge theory variables** [Gambini & Trias 81], [Jacobson, Rovelli, Smolin 88]

- Magnet. dof.: Holonomy (Wilson – Loop)

$$A(e) = \mathcal{P} \exp\left(\int_e A\right)$$

- Electr. dof: flux

$$E_j(S) = \int_S \epsilon_{abc} E_j^a dx^b \wedge dx^c$$

- Poisson – brackets:

$$\{E_j(S), A(e)\} = G A(e_1) \tau_j A(e_2); \quad e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$

# Physical Hilbert Space

**Lattice – inspired canon. gauge theory variables** [Gambini & Trias 81], [Jacobson, Rovelli, Smolin 88]

- Magnet. dof.: Holonomy (Wilson – Loop)

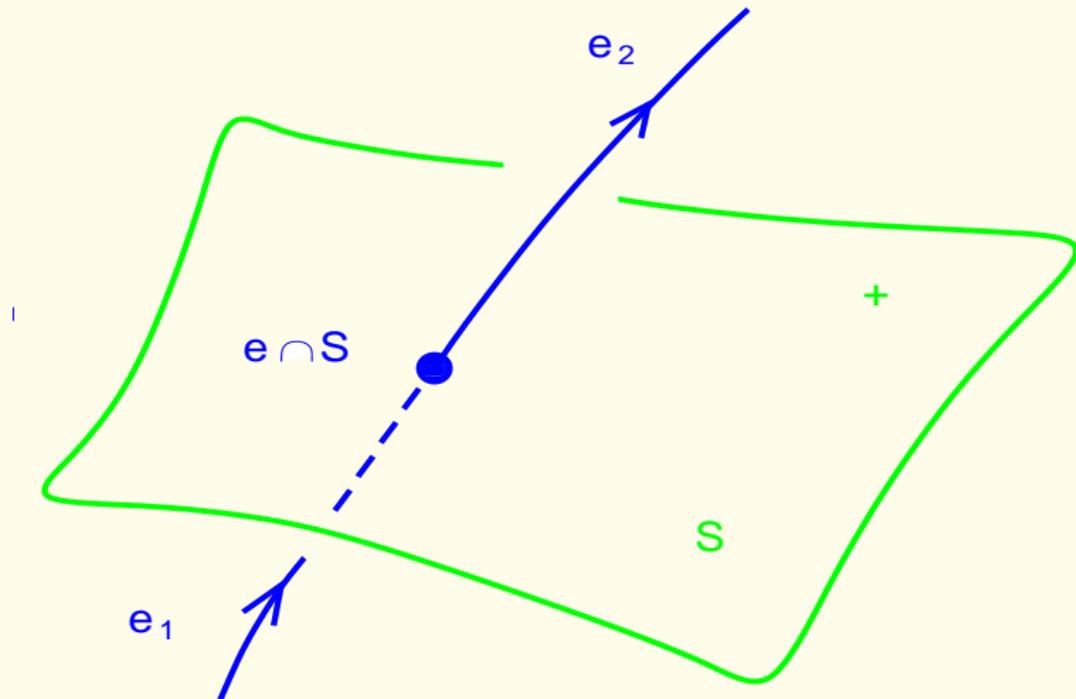
$$A(e) = \mathcal{P} \exp\left(\int_e A\right)$$

- Electr. dof: flux

$$E_j(S) = \int_S \epsilon_{abc} E_j^a dx^b \wedge dx^c$$

- Poisson – brackets:

$$\{E_j(S), A(e)\} = G A(e_1) \tau_j A(e_2); \quad e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$



## Lattice – inspired gauge theory variables [Gambini & Trias 81], [Jacobson, Rovelli, Smolin 88]

- Magnet. dof.: Holonomy (Wilson – Loop)

$$A(e) = \mathcal{P} \exp\left(\int_e A\right)$$

- Electr. dof: flux

$$E_f(S) = \int_S \epsilon_{abc} E_j^a dx^b \wedge dx^c$$

- Poisson – brackets:

$$\{E_j(S), A(e)\} = G A(e_1) \tau_j A(e_2); \quad e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$

- Reality conditions:

$$\overline{A(e)} = [A(e^{-1})]^T, \quad \overline{E_j(S)} = E_j(S)$$

- Defines abstract Poisson\* – algebra  $\mathfrak{A}_{\text{phys}}$ .
- Bundle automorphisms  $\mathfrak{G} \cong \mathcal{G} \times \text{Diff}(\Sigma)$  act by Poisson automorphisms on  $\mathfrak{A}_{\text{phys}}$  e.g.  $\alpha_g = \exp(\int \lambda^i c_i, \cdot)$ ,  $g = \exp(\lambda^i \tau_i)$

$$\alpha_g(A(e)) = g(b(e)) A(e) g(f(e))^{-1}, \quad \alpha_\varphi(A(e)) = A(\varphi(e))$$

## Lattice – inspired gauge theory variables [Gambini & Trias 81], [Jacobson, Rovelli, Smolin 88]

- Magnet. dof.: Holonomy (Wilson – Loop)

$$A(e) = \mathcal{P} \exp\left(\int_e A\right)$$

- Electr. dof: flux

$$E_f(S) = \int_S \epsilon_{abc} E_j^a dx^b \wedge dx^c$$

- Poisson – brackets:

$$\{E_j(S), A(e)\} = G A(e_1) \tau_j A(e_2); \quad e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$

- Reality conditions:

$$\overline{A(e)} = [A(e^{-1})]^T, \quad \overline{E_j(S)} = E_j(S)$$

- Defines abstract Poisson\* – algebra  $\mathfrak{A}_{\text{phys}}$ .
- Bundle automorphisms  $\mathfrak{G} \cong \mathcal{G} \times \text{Diff}(\Sigma)$  act by Poisson automorphisms on  $\mathfrak{A}_{\text{phys}}$  e.g.  $\alpha_g = \exp(\int \lambda^j c_j, \cdot)$ ,  $g = \exp(\lambda^j \tau_j)$

$$\alpha_g(A(e)) = g(b(e)) A(e) g(f(e))^{-1}, \quad \alpha_\varphi(A(e)) = A(\varphi(e))$$

## Lattice – inspired gauge theory variables [Gambini & Trias 81], [Jacobson, Rovelli, Smolin 88]

- Magnet. dof.: Holonomy (Wilson – Loop)

$$A(e) = \mathcal{P} \exp\left(\int_e A\right)$$

- Electr. dof: flux

$$E_f(S) = \int_S \epsilon_{abc} E_j^a dx^b \wedge dx^c$$

- Poisson – brackets:

$$\{E_j(S), A(e)\} = G A(e_1) \tau_j A(e_2); \quad e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$

- Reality conditions:

$$\overline{A(e)} = [A(e^{-1})]^T, \quad \overline{E_j(S)} = E_j(S)$$

- Defines abstract Poisson\* – algebra  $\mathfrak{A}_{\text{phys}}$ .
- Bundle automorphisms  $\mathfrak{G} \cong \mathcal{G} \times \text{Diff}(\Sigma)$  act by Poisson automorphisms on  $\mathfrak{A}_{\text{phys}}$  e.g.  $\alpha_g = \exp(\{ \int \lambda^j c_j, \cdot \})$ ,  $g = \exp(\lambda^j \tau_j)$

$$\alpha_g(A(e)) = g(b(e)) A(e) g(f(e))^{-1}, \quad \alpha_\varphi(A(e)) = A(\varphi(e))$$

- **HS Repr.:** In QFT no Stone – von Neumann Theorem!!!

**THEOREM** [Ashtekar, Isham, Lewandowski 92-93], [Sahlmann 02], [L., Okolow, S., T.T. 03-05], [Fleischhack 04]

Diff( $\Sigma$ ) inv. states on hol. – flux algebra  $\mathfrak{A}_{\text{phys}}$  unique.

- wave functions of  $\mathcal{H}_{\text{phys}}$

$$\psi(A) = \psi_\gamma(A(\mathbf{e}_1), \dots, A(\mathbf{e}_N)), \quad \psi_\gamma : \text{SU}(2)^N \rightarrow \mathbb{C}$$

- Holonomy = multiplication – operator

$$[\widehat{A(\mathbf{e})} \psi](A) := A(\mathbf{e}) \psi(A)$$

- Flux = derivative – operator

$$[\widehat{E_j(S)} \psi](A) := i\hbar \{E_j(S), \psi(A)\}$$

- Scalar product

$$\langle \psi, \psi' \rangle := \int_{\text{SU}(2)^N} d\mu_H(h_1) \dots d\mu_H(h_N) \overline{\psi_\gamma(h_1, \dots, h_N)} \psi'_\gamma(h_1, \dots, h_N)$$

- **HS Reprs.:** In QFT no Stone – von Neumann Theorem!!!

**Theorem** [Ashtekar, Isham, Lewandowski 92-93], [Sahlmann 02], [L., Okolow, S., T.T. 03-05], [Fleischhack 04]

Diff( $\Sigma$ ) inv. states on hol. – flux algebra  $\mathfrak{A}_{\text{phys}}$  unique.

- wave functions of  $\mathcal{H}_{\text{phys}}$

$$\psi(A) = \psi_\gamma(A(e_1), \dots, A(e_N)), \quad \psi_\gamma : \text{SU}(2)^N \rightarrow \mathbb{C}$$

- Holonomy = multiplication – operator

$$[\widehat{A(e)} \psi](A) := A(e) \psi(A)$$

- Flux = derivative – operator

$$[\widehat{E_j(S)} \psi](A) := i\hbar \{E_j(S), \psi(A)\}$$

- Scalar product

$$\langle \psi, \psi' \rangle := \int_{\text{SU}(2)^N} d\mu_H(h_1) \dots d\mu_H(h_N) \overline{\psi_\gamma(h_1, \dots, h_N)} \psi'_\gamma(h_1, \dots, h_N)$$

- **HS Reprs.:** In QFT no Stone – von Neumann Theorem!!!

**Theorem** [Ashtekar, Isham, Lewandowski 92-93], [Sahlmann 02], [L., Okolow, S., T.T. 03-05], [Fleischhack 04]

Diff( $\Sigma$ ) inv. states on hol. – flux algebra  $\mathfrak{A}_{\text{phys}}$  unique.

- wave functions of  $\mathcal{H}_{\text{phys}}$

$$\psi(A) = \psi_\gamma(A(\mathbf{e}_1), \dots, A(\mathbf{e}_N)), \quad \psi_\gamma : \text{SU}(2)^N \rightarrow \mathbb{C}$$

- Holonomy = multiplication – operator

$$[\widehat{A(\mathbf{e})} \psi](A) := A(\mathbf{e}) \psi(A)$$

- Flux = derivative – operator

$$[\widehat{E_j(S)} \psi](A) := i\hbar \{E_j(S), \psi(A)\}$$

- Scalar product

$$\langle \psi, \psi' \rangle := \int_{\text{SU}(2)^N} d\mu_H(h_1) \dots d\mu_H(h_N) \overline{\psi_\gamma(h_1, \dots, h_N)} \psi'_\gamma(h_1, \dots, h_N)$$

- **HS Reprs.:** In QFT no Stone – von Neumann Theorem!!!

**Theorem** [Ashtekar, Isham, Lewandowski 92-93], [Sahlmann 02], [L., Okolow, S., T.T. 03-05], [Fleischhack 04]

Diff( $\Sigma$ ) inv. states on hol. – flux algebra  $\mathfrak{A}_{\text{phys}}$  unique.

- wave functions of  $\mathcal{H}_{\text{phys}}$

$$\psi(A) = \psi_\gamma(A(\mathbf{e}_1), \dots, A(\mathbf{e}_N)), \quad \psi_\gamma : \text{SU}(2)^N \rightarrow \mathbb{C}$$

- Holonomy = multiplication – operator

$$[\widehat{A(\mathbf{e})} \psi](A) := A(\mathbf{e}) \psi(A)$$

- Flux = derivative – operator

$$[\widehat{E_j(S)} \psi](A) := i\hbar \{E_j(S), \psi(A)\}$$

- Scalar product

$$\langle \psi, \psi' \rangle := \int_{\text{SU}(2)^N} d\mu_H(h_1) \dots d\mu_H(h_N) \overline{\psi_\gamma(h_1, \dots, h_N)} \psi'_\gamma(h_1, \dots, h_N)$$

- **HS Reprs.:** In QFT no Stone – von Neumann Theorem!!!

**Theorem** [Ashtekar, Isham, Lewandowski 92-93], [Sahlmann 02], [L., Okolow, S., T.T. 03-05], [Fleischhack 04]

Diff( $\Sigma$ ) inv. states on hol. – flux algebra  $\mathfrak{A}_{\text{phys}}$  unique.

- wave functions of  $\mathcal{H}_{\text{phys}}$

$$\psi(A) = \psi_\gamma(A(\mathbf{e}_1), \dots, A(\mathbf{e}_N)), \quad \psi_\gamma : \text{SU}(2)^N \rightarrow \mathbb{C}$$

- Holonomy = multiplication – operator

$$[\widehat{A(\mathbf{e})} \psi](A) := A(\mathbf{e}) \psi(A)$$

- Flux = derivative – operator

$$[\widehat{E_j(\mathbf{S})} \psi](A) := i\hbar \{E_j(\mathbf{S}), \psi(A)\}$$

- Scalar product

$$\langle \psi, \psi' \rangle := \int_{\text{SU}(2)^N} d\mu_H(h_1) \dots d\mu_H(h_N) \overline{\psi_\gamma(h_1, \dots, h_N)} \psi'_\gamma(h_1, \dots, h_N)$$

- **HS Reprs.:** In QFT no Stone – von Neumann Theorem!!!

**Theorem** [Ashtekar, Isham, Lewandowski 92-93], [Sahlmann 02], [L., Okolow, S., T.T. 03-05], [Fleischhack 04]

Diff( $\Sigma$ ) inv. states on hol. – flux algebra  $\mathfrak{A}_{\text{phys}}$  unique.

- wave functions of  $\mathcal{H}_{\text{phys}}$

$$\psi(A) = \psi_{\gamma}(A(\mathbf{e}_1), \dots, A(\mathbf{e}_N)), \quad \psi_{\gamma} : \text{SU}(2)^N \rightarrow \mathbb{C}$$

- Holonomy = multiplication – operator

$$[\widehat{A(\mathbf{e})} \psi](A) := A(\mathbf{e}) \psi(A)$$

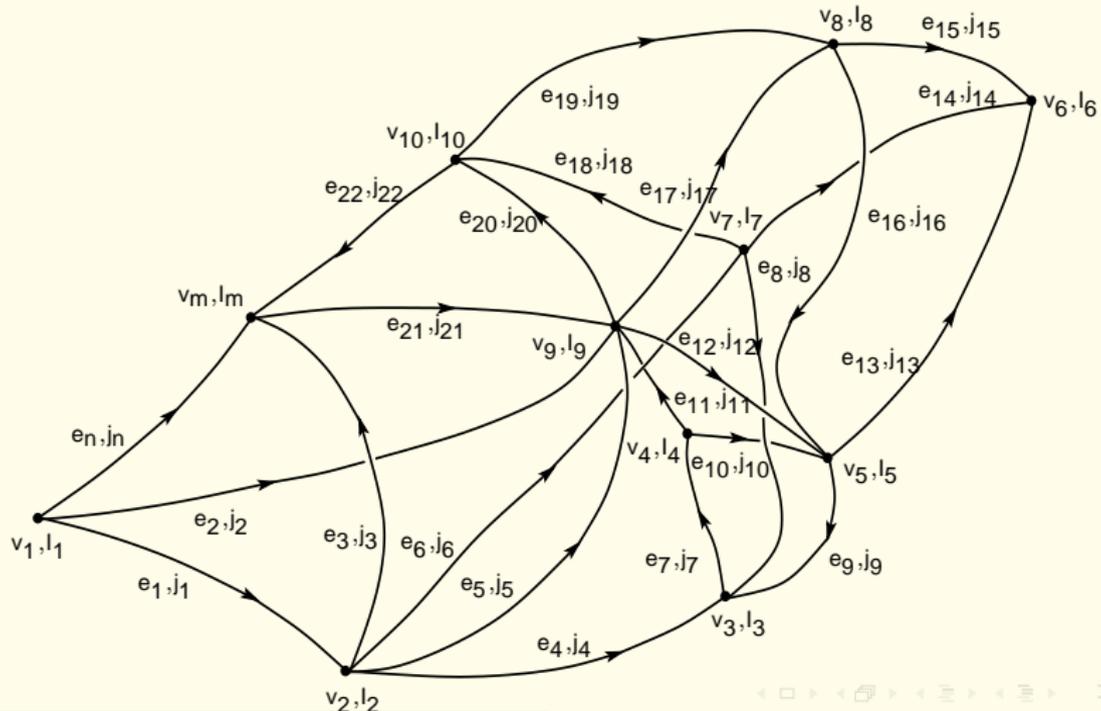
- Flux = derivative – operator

$$[\widehat{E_j(\mathbf{S})} \psi](A) := i\hbar \{E_j(\mathbf{S}), \psi(A)\}$$

- Scalar product

$$\langle \psi, \psi' \rangle := \int_{\text{SU}(2)^N} d\mu_{\text{H}}(h_1) \dots d\mu_{\text{H}}(h_N) \overline{\psi_{\gamma}(h_1, \dots, h_N)} \psi'_{\gamma}(h_1, \dots, h_N)$$

## Spin Network ONB $T_{\gamma,j,l}$



## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners I
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{T_{\gamma,j,l}; j \neq 0, l\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)T_{\gamma,j,l} := T_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$
2. Compute  $\langle \psi_{\gamma}, H \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners I
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, l\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$
2. Compute  $\langle \psi_{\gamma}, H \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners I
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, l\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$
2. Compute  $\langle \psi_{\gamma}, H \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners  $\mathbb{I}$
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, \mathbb{I}\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$
2. Compute  $\langle \psi_{\gamma}, H \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners I
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, l\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$

2. Compute  $\langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners  $\mathbb{I}$
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, \mathbb{I}\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma}, \forall \gamma$

2. Compute  $\langle \psi_{\gamma}, \hat{H}_{\gamma} \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle + f. \text{ semiclassical } \psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners I
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, l\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$

2. Compute  $\langle \psi_{\gamma}, H \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

## Does rep. support $\hat{H}$ with correct semiclassical limit?

- Gauss constraint solved by restriction of intertwiners I
- $\mathcal{H}_{\text{phys}}$  not separable

$$\mathcal{H}_{\text{phys}} = \bigoplus_{\gamma} \mathcal{H}_{\gamma}, \quad \mathcal{H}_{\gamma} = \overline{\text{span}\{\mathbb{T}_{\gamma,j,l}; j \neq 0, l\}}$$

- $\text{Diff}(\Sigma)$  does not downsize it since symmetry group, not gauge group
- Unitary representation  $U(\varphi)\mathbb{T}_{\gamma,j,l} := \mathbb{T}_{\varphi(\gamma),j,l}$
- If  $U(\varphi) F U(\varphi)^{-1} = F$  (e.g.  $F = H$ ; all operationally defined observables) then “superselection” (subgraph preservation)

$$F \mathcal{H}_{\gamma} \subset \mathcal{H}_{\gamma} \Rightarrow F = \bigoplus_{\gamma} F_{\gamma}$$

- This imposes strong constraints on regularisation of  $\hat{H}$  and removes most ambiguities usually encountered for  $\hat{C}$  !
- Task:

1. Construct  $\hat{H}_{\gamma} \forall \gamma$
2. Compute  $\langle \psi_{\gamma}, H \psi_{\gamma} \rangle = \langle \psi_{\gamma}, H_{\gamma} \psi_{\gamma} \rangle$  f. semiclass.  $\psi_{\gamma}$

# Physical coherent states

- Choose cell complex  $\gamma^*$ , dual graph  $\gamma$  s.t.  $e \leftrightarrow S_e$
- Choose classical field configuration  $(A_0(x), E_0(x))$ , compute

$$g_e := \exp(i\tau_j E_0^j(S_e)) A_0(e) \in G^C$$

- Define [Hall 90's], [Sahlmann, T.T., Winkler 00's]

$$\psi_{A_0, E_0} := \otimes_e \psi_e, \quad \psi_e(h_e) := \sum_{\pi} \dim(\pi) e^{-\tau_e \lambda_{\pi}} \chi_{\pi}(g_e h_e^{-1})$$

- Minimal uncertainty states, that is,  $\forall e \in E(\gamma)$

$$\langle \psi_{A_0, E_0}, \widehat{A}(e) \psi_{A_0, E_0} \rangle = A_0(e), \quad \langle \psi_{A_0, E_0}, \widehat{E}_j(S_e) \psi_{A_0, E_0} \rangle = E_{j0}(S_e)$$

$$\langle \widehat{\Delta A}(e) \rangle \langle \widehat{\Delta E}_j(S_e) \rangle = \frac{1}{2} | \langle \widehat{[A}(e), E_j(S_e)] \rangle |$$

## Physical coherent states

- Choose cell complex  $\gamma^*$ , dual graph  $\gamma$  s.t.  $e \leftrightarrow S_e$
- Choose classical field configuration  $(A_0(x), E_0(x))$ , compute

$$g_e := \exp(i\tau_j E_0^j(S_e)) A_0(e) \in G^C$$

- Define [Hall 90's], [Sahlmann, T.T., Winkler 00's]

$$\psi_{A_0, E_0} := \otimes_e \psi_e, \quad \psi_e(h_e) := \sum_{\pi} \dim(\pi) e^{-\tau_e \lambda_{\pi}} \chi_{\pi}(g_e h_e^{-1})$$

- Minimal uncertainty states, that is,  $\forall e \in E(\gamma)$

$$\langle \psi_{A_0, E_0}, \widehat{A}(e) \psi_{A_0, E_0} \rangle = A_0(e), \quad \langle \psi_{A_0, E_0}, \widehat{E}_j(S_e) \psi_{A_0, E_0} \rangle = E_{j0}(S_e)$$

$$\langle \Delta \widehat{A}(e) \rangle \langle \Delta \widehat{E}_j(S_e) \rangle = \frac{1}{2} | \langle \widehat{[A(e), E_j(S_e)]} \rangle |$$

## Physical coherent states

- Choose cell complex  $\gamma^*$ , dual graph  $\gamma$  s.t.  $e \leftrightarrow S_e$
- Choose classical field configuration  $(A_0(x), E_0(x))$ , compute

$$g_e := \exp(i\tau_j E_0^j(S_e)) A_0(e) \in G^{\mathbb{C}}$$

- Define [Hall 90's], [Sahlmann, T.T., Winkler 00's]

$$\psi_{A_0, E_0} := \otimes_e \psi_e, \quad \psi_e(\mathbf{h}_e) := \sum_{\pi} \dim(\pi) e^{-\tau_e \lambda_{\pi}} \chi_{\pi}(g_e \mathbf{h}_e^{-1})$$

- Minimal uncertainty states, that is,  $\forall e \in E(\gamma)$

$$\langle \psi_{A_0, E_0}, \widehat{A}(e) \psi_{A_0, E_0} \rangle = A_0(e), \quad \langle \psi_{A_0, E_0}, \widehat{E}_j(S_e) \psi_{A_0, E_0} \rangle = E_{j0}(S_e)$$

$$\langle \Delta \widehat{A}(e) \rangle \langle \Delta \widehat{E}_j(S_e) \rangle = \frac{1}{2} | \langle \widehat{[A(e), E_j(S_e)]} \rangle |$$

## Physical coherent states

- Choose cell complex  $\gamma^*$ , dual graph  $\gamma$  s.t.  $e \leftrightarrow S_e$
- Choose classical field configuration  $(A_0(x), E_0(x))$ , compute

$$g_e := \exp(i\tau_j E_0^j(S_e)) A_0(e) \in G^{\mathbb{C}}$$

- Define [Hall 90's], [Sahlmann, T.T., Winkler 00's]

$$\psi_{A_0, E_0} := \otimes_e \psi_e, \quad \psi_e(h_e) := \sum_{\pi} \dim(\pi) e^{-\tau_e \lambda_{\pi}} \chi_{\pi}(g_e h_e^{-1})$$

- Minimal uncertainty states, that is,  $\forall e \in E(\gamma)$

1.

$$\langle \psi_{A_0, E_0}, \widehat{A(e)} \psi_{A_0, E_0} \rangle = A_0(e), \quad \langle \psi_{A_0, E_0}, \widehat{E_j(S_e)} \psi_{A_0, E_0} \rangle = E_{j0}(S_e)$$

2.

$$\langle \widehat{\Delta A(e)} \rangle \langle \widehat{\Delta E_j(S_e)} \rangle = \frac{1}{2} | \langle [\widehat{A(e)}, \widehat{E_j(S_e)}] \rangle |$$

## Physical coherent states

- Choose cell complex  $\gamma^*$ , dual graph  $\gamma$  s.t.  $e \leftrightarrow S_e$
- Choose classical field configuration  $(A_0(x), E_0(x))$ , compute

$$g_e := \exp(i\tau_j E_0^j(S_e)) A_0(e) \in G^{\mathbb{C}}$$

- Define [Hall 90's], [Sahlmann, T.T., Winkler 00's]

$$\psi_{A_0, E_0} := \otimes_e \psi_e, \quad \psi_e(h_e) := \sum_{\pi} \dim(\pi) e^{-\tau_e \lambda_{\pi}} \chi_{\pi}(g_e h_e^{-1})$$

- Minimal uncertainty states, that is,  $\forall e \in E(\gamma)$

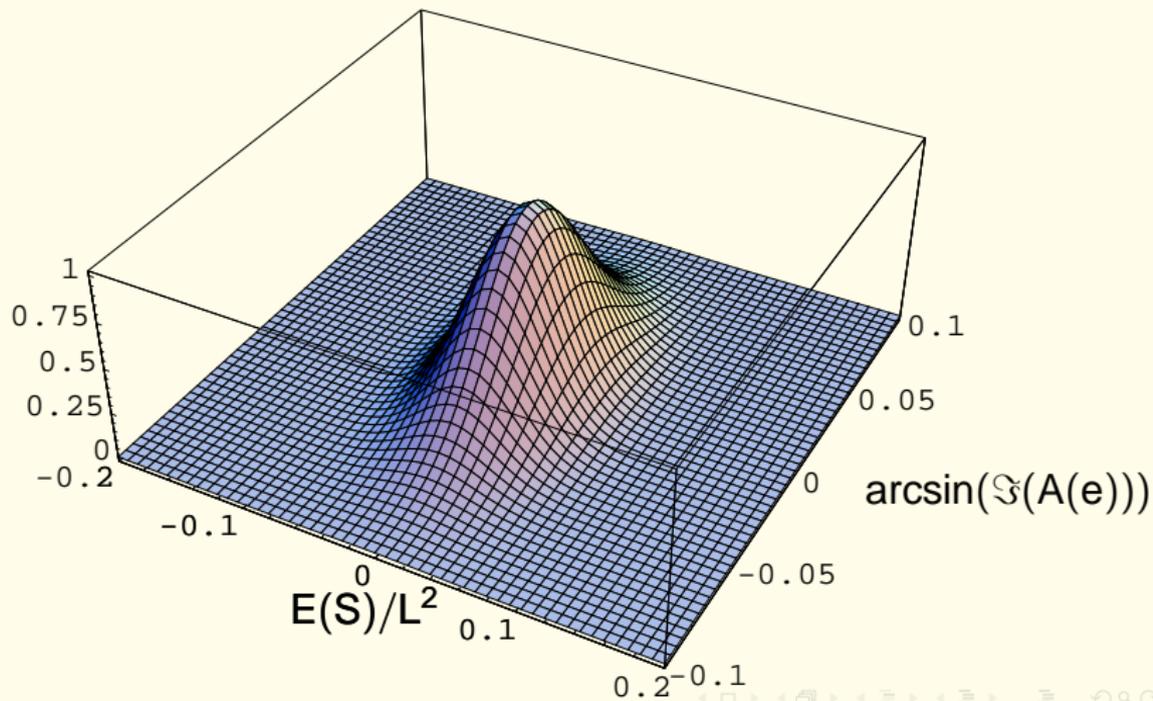
1.

$$\langle \psi_{A_0, E_0}, \widehat{A}(e) \psi_{A_0, E_0} \rangle = A_0(e), \quad \langle \psi_{A_0, E_0}, \widehat{E}_j(S_e) \psi_{A_0, E_0} \rangle = E_{j0}(S_e)$$

2.

$$\langle \widehat{\Delta A}(e) \rangle \langle \widehat{\Delta E}_j(S_e) \rangle = \frac{1}{2} | \langle [\widehat{A}(e), \widehat{E}_j(S_e)] \rangle |$$

## Overlap Function



## Remarks:

- Notice:  $\Sigma$  just differential manifold, no Riemannian space!
- No a priori meaning to how densely  $\gamma$  embedded
- In particular, final operator  $\hat{H}$  cannot depend on short distance regulator used at intermediate stages of construction
- Notice: Operator family  $(\hat{H}_\gamma)$  defines **Continuum operator**
- Expect that good semiclassical states depend on graphs which are very densely embedded wrt background metric to be approximated
- Choose graph to be countably infinite (for compact  $\Sigma$  large finite graph sufficient)

## Remarks:

- Notice:  $\Sigma$  just differential manifold, no Riemannian space!
- No a priori meaning to how densely  $\gamma$  embedded
- In particular, final operator  $\hat{H}$  cannot depend on short distance regulator used at intermediate stages of construction
- Notice: Operator family  $(\hat{H}_\gamma)$  defines **Continuum operator**
- Expect that good semiclassical states depend on graphs which are very densely embedded wrt background metric to be approximated
- Choose graph to be countably infinite (for compact  $\Sigma$  large finite graph sufficient)

## Remarks:

- Notice:  $\Sigma$  just differential manifold, no Riemannian space!
- No a priori meaning to how densely  $\gamma$  embedded
- In particular, final operator  $\hat{H}$  cannot depend on short distance regulator used at intermediate stages of construction
- Notice: Operator family  $(\hat{H}_\gamma)$  defines **Continuum operator**
- Expect that good semiclassical states depend on graphs which are very densely embedded wrt background metric to be approximated
- Choose graph to be countably infinite (for compact  $\Sigma$  large finite graph sufficient)

## Remarks:

- Notice:  $\Sigma$  just differential manifold, no Riemannian space!
- No a priori meaning to how densely  $\gamma$  embedded
- In particular, final operator  $\hat{H}$  cannot depend on short distance regulator used at intermediate stages of construction
- Notice: Operator family  $(\hat{H}_\gamma)$  defines **Continuum operator**
- Expect that good semiclassical states depend on graphs which are very densely embedded wrt background metric to be approximated
- Choose graph to be countably infinite (for compact  $\Sigma$  large finite graph sufficient)

## Remarks:

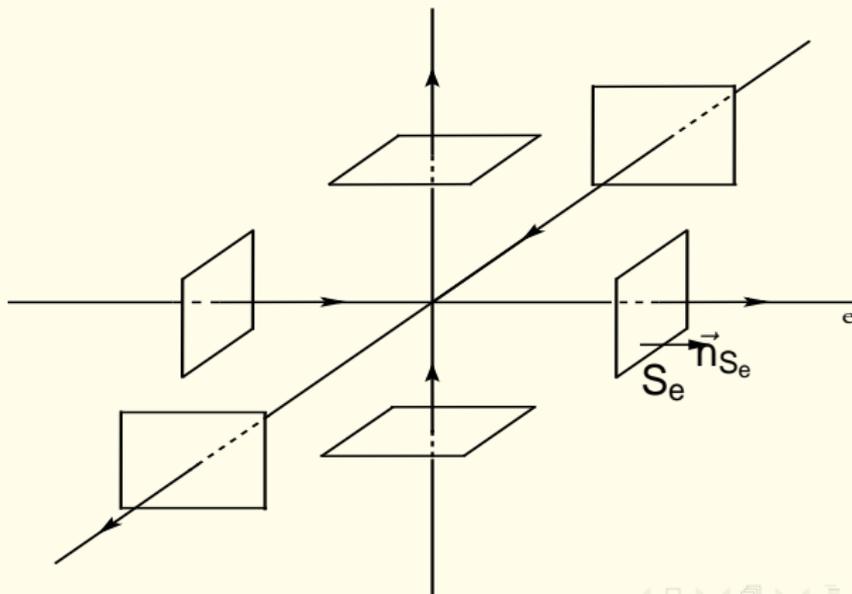
- Notice:  $\Sigma$  just differential manifold, no Riemannian space!
- No a priori meaning to how densely  $\gamma$  embedded
- In particular, final operator  $\hat{H}$  cannot depend on short distance regulator used at intermediate stages of construction
- Notice: Operator family  $(\hat{H}_\gamma)$  defines **Continuum operator**
- Expect that good semiclassical states depend on graphs which are very densely embedded wrt background metric to be approximated
- Choose graph to be countably infinite (for compact  $\Sigma$  large finite graph sufficient)

## Remarks:

- Notice:  $\Sigma$  just differential manifold, no Riemannian space!
- No a priori meaning to how densely  $\gamma$  embedded
- In particular, final operator  $\hat{H}$  cannot depend on short distance regulator used at intermediate stages of construction
- Notice: Operator family  $(\hat{H}_\gamma)$  defines **Continuum operator**
- Expect that good semiclassical states depend on graphs which are very densely embedded wrt background metric to be approximated
- Choose graph to be countably infinite (for compact  $\Sigma$  large finite graph sufficient)

# Physical Hamiltonian

## Example: Cubic graph



## Comparison with YM theory on cubic lattice

- Yang – Mills on  $(\mathbb{R}^4, \eta)$  [Kogut & Susskind 74]

$$H_\gamma = \frac{\hbar}{2g^2} \epsilon \sum_{v \in V(\gamma)} \sum_{a=1}^3 \text{Tr} \left( E(S_v^a)^2 + [2 - A(\alpha_v^a) - A(\alpha_v^a)^{-1}] \right)$$

- Gravity on  $\mathbb{R} \times \Sigma$  [T.T. 96 – 05, Giesel & T.T. 06]

$$H_\gamma = \frac{\hbar}{\ell_P^4} \sum_{v \in V(\gamma)} \sqrt{\left| \sum_{\mu=0}^3 \eta^{\mu\mu} \left[ \sum_{a=1}^3 \text{Tr} (\tau_\mu A(\alpha_v^a) A(e_v^a) [A(e_v^a)^{-1}, V_v]) \right]^2 \right|}$$

- Volume operator

$$V_v = \sqrt{|\epsilon_{abc} \text{Tr} (E(S_v^a) E(S_v^b) E(S_v^c))|}$$

- Lattice spacing  $\epsilon$  disappears, automat. UV finite.
- In a precise sense:  $\epsilon$  replaced by  $\ell_P$

## Comparison with YM theory on cubic lattice

- Yang – Mills on  $(\mathbb{R}^4, \eta)$  [Kogut & Susskind 74]

$$H_\gamma = \frac{\hbar}{2g^2} \epsilon \sum_{v \in V(\gamma)} \sum_{a=1}^3 \text{Tr} \left( E(S_v^a)^2 + [2 - A(\alpha_v^a) - A(\alpha_v^a)^{-1}] \right)$$

- Gravity on  $\mathbb{R} \times \Sigma$  [T.T. 96 – 05, Giesel & T.T. 06]

$$H_\gamma = \frac{\hbar}{\ell_P^4} \sum_{v \in V(\gamma)} \sqrt{\left| \sum_{\mu=0}^3 \eta^{\mu\mu} \left[ \sum_{a=1}^3 \text{Tr} (\tau_\mu A(\alpha_v^a) A(e_v^a) [A(e_v^a)^{-1}, V_v]) \right]^2 \right|}$$

- Volume operator

$$V_v = \sqrt{|\epsilon_{abc} \text{Tr} (E(S_v^a) E(S_v^b) E(S_v^c))|}$$

- Lattice spacing  $\epsilon$  disappears, automat. UV finite.
- In a precise sense:  $\epsilon$  replaced by  $\ell_P$

## Comparison with YM theory on cubic lattice

- Yang – Mills on  $(\mathbb{R}^4, \eta)$  [Kogut & Susskind 74]

$$H_\gamma = \frac{\hbar}{2g^2} \epsilon \sum_{v \in V(\gamma)} \sum_{a=1}^3 \text{Tr} \left( \mathbf{E}(\mathbf{S}_v^a)^2 + [2 - A(\alpha_v^a) - A(\alpha_v^a)^{-1}] \right)$$

- Gravity on  $\mathbb{R} \times \Sigma$  [T.T. 96 – 05, Giesel & T.T. 06]

$$H_\gamma = \frac{\hbar}{\ell_P^4} \sum_{v \in V(\gamma)} \sqrt{\left| \sum_{\mu=0}^3 \eta^{\mu\mu} \left[ \sum_{a=1}^3 \text{Tr} (\tau_\mu A(\alpha_v^a) A(\mathbf{e}_v^a) [A(\mathbf{e}_v^a)^{-1}, V_v]) \right]^2 \right|}$$

- Volume operator

$$V_v = \sqrt{|\epsilon_{abc} \text{Tr} (\mathbf{E}(\mathbf{S}_v^a) \mathbf{E}(\mathbf{S}_v^b) \mathbf{E}(\mathbf{S}_v^c))|}$$

- Lattice spacing  $\epsilon$  disappears, automat. UV finite.
- In a precise sense:  $\epsilon$  replaced by  $\ell_P$

## Comparison with YM theory on cubic lattice

- Yang – Mills on  $(\mathbb{R}^4, \eta)$  [Kogut & Susskind 74]

$$H_\gamma = \frac{\hbar}{2g^2} \epsilon \sum_{v \in V(\gamma)} \sum_{a=1}^3 \text{Tr} \left( \mathbf{E}(\mathbf{S}_v^a)^2 + [2 - A(\alpha_v^a) - A(\alpha_v^a)^{-1}] \right)$$

- Gravity on  $\mathbb{R} \times \Sigma$  [T.T. 96 – 05, Giesel & T.T. 06]

$$H_\gamma = \frac{\hbar}{\ell_P^4} \sum_{v \in V(\gamma)} \sqrt{\left| \sum_{\mu=0}^3 \eta^{\mu\mu} \left[ \sum_{a=1}^3 \text{Tr} (\tau_\mu A(\alpha_v^a) A(\mathbf{e}_v^a) [A(\mathbf{e}_v^a)^{-1}, V_v]) \right]^2 \right|}$$

- Volume operator

$$V_v = \sqrt{|\epsilon_{abc} \text{Tr} (\mathbf{E}(\mathbf{S}_v^a) \mathbf{E}(\mathbf{S}_v^b) \mathbf{E}(\mathbf{S}_v^c))|}$$

- Lattice spacing  $\epsilon$  disappears, automat. UV finite.
- In a precise sense:  $\epsilon$  replaced by  $\ell_P$

## Comparison with YM theory on cubic lattice

- Yang – Mills on  $(\mathbb{R}^4, \eta)$  [Kogut & Susskind 74]

$$H_\gamma = \frac{\hbar}{2g^2} \epsilon \sum_{v \in V(\gamma)} \sum_{a=1}^3 \text{Tr} \left( \mathbf{E}(\mathbf{S}_v^a)^2 + [2 - A(\alpha_v^a) - A(\alpha_v^a)^{-1}] \right)$$

- Gravity on  $\mathbb{R} \times \Sigma$  [T.T. 96 – 05, Giesel & T.T. 06]

$$H_\gamma = \frac{\hbar}{\ell_P^4} \sum_{v \in V(\gamma)} \sqrt{\left| \sum_{\mu=0}^3 \eta^{\mu\mu} \left[ \sum_{a=1}^3 \text{Tr} (\tau_\mu A(\alpha_v^a) A(\mathbf{e}_v^a) [A(\mathbf{e}_v^a)^{-1}, V_v]) \right]^2 \right|}$$

- Volume operator

$$V_v = \sqrt{|\epsilon_{abc} \text{Tr} (\mathbf{E}(\mathbf{S}_v^a) \mathbf{E}(\mathbf{S}_v^b) \mathbf{E}(\mathbf{S}_v^c))|}$$

- Lattice spacing  $\epsilon$  disappears, automat. UV finite.
- In a precise sense:  $\epsilon$  replaced by  $\ell_P$

# Semiclassical Limit

**Theorem** [Giesel & T.T. 06] For any  $(A_0, E_0)$ , suff. large  $\gamma$

1. Exp. Value  $\langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle = H(A_0, E_0) + O(\hbar)$
2. Fluctuation  $\langle \psi_{A_0, E_0}, \hat{H}^2 \psi_{A_0, E_0} \rangle - \langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle^2 = O(\hbar)$

**Corollary**

Quantum Hamiltonian correctly implemented

For sufficiently small  $\hbar$

$$e^{i\hat{H}/\hbar} \psi_{A_0, E_0} \approx \psi_{A_0(\tau), E_0(\tau)}$$

# Semiclassical Limit

**Theorem** [Giesel & T.T. 06] For any  $(A_0, E_0)$ , suff. large  $\gamma$

1. Exp. Value  $\langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle = H(A_0, E_0) + O(\hbar)$
2. Fluctuation  $\langle \psi_{A_0, E_0}, \hat{H}^2 \psi_{A_0, E_0} \rangle - \langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle^2 = O(\hbar)$

**Corollary**

Quantum Hamiltonian correctly implemented

For sufficiently small  $\tau$

$$e^{+\tau \hat{H} / \hbar} \psi_{A_0, E_0} \approx \psi_{A_0(\tau), E_0(\tau)}$$

## Semiclassical Limit

**Theorem** [Giesel & T.T. 06] For any  $(A_0, E_0)$ , suff. large  $\gamma$

1. Exp. Value  $\langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle = \mathbf{H}(A_0, E_0) + \mathcal{O}(\hbar)$
2. Fluctuation  $\langle \psi_{A_0, E_0}, \hat{H}^2 \psi_{A_0, E_0} \rangle - \langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle^2 = \mathcal{O}(\hbar)$

**Corollary**

Quantum Hamiltonian correctly implemented

For sufficiently small  $\hbar$

$$e^{+i\hat{H}/\hbar} \psi_{A_0, E_0} \approx \psi_{A_0(\tau), E_0(\tau)}$$

## Semiclassical Limit

Theorem [Giesel & T.T. 06] For any  $(A_0, E_0)$ , suff. large  $\gamma$

1. Exp. Value  $\langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle = \mathbf{H}(A_0, E_0) + \mathcal{O}(\hbar)$
2. Fluctuation  $\langle \psi_{A_0, E_0}, \hat{H}^2 \psi_{A_0, E_0} \rangle - \langle \psi_{A_0, E_0}, \hat{H} \psi_{A_0, E_0} \rangle^2 = \mathcal{O}(\hbar)$

Corollary

- i. Quantum Hamiltonian correctly implemented
- ii. For sufficiently small  $\tau$

$$e^{i\tau\hat{H}/\hbar} \psi_{A_0, E_0} \approx \psi_{A_0(\tau), E_0(\tau)}$$

## Semiclassical Limit

Theorem [Giesel & T.T. 06] For any  $(A_0, E_0)$ , suff. large  $\gamma$

1. Exp. Value  $\langle \psi_{A_0, E_0}, \hat{\mathbf{H}} \psi_{A_0, E_0} \rangle = \mathbf{H}(A_0, E_0) + \mathcal{O}(\hbar)$
2. Fluctuation  $\langle \psi_{A_0, E_0}, \hat{\mathbf{H}}^2 \psi_{A_0, E_0} \rangle - \langle \psi_{A_0, E_0}, \hat{\mathbf{H}} \psi_{A_0, E_0} \rangle^2 = \mathcal{O}(\hbar)$

Corollary

- i. Quantum Hamiltonian correctly implemented
- ii. For sufficiently small  $\tau$

$$e^{i\tau\hat{\mathbf{H}}/\hbar} \psi_{A_0, E_0} \approx \psi_{A_0(\tau), E_0(\tau)}$$

# Summary

- LQG dynamically severely constrained (uniqueness result)
- correct semiclassical limit of  $\hat{H}$  established
- Final picture equivalent to background independent, Hamiltonian “floating” lattice gauge theory

# Summary

- LQG dynamically severely constrained (uniqueness result)
- correct semiclassical limit of  $\hat{H}$  established
- Final picture equivalent to background independent, Hamiltonian “floating” lattice gauge theory

# Summary

- LQG dynamically severely constrained (uniqueness result)
- correct semiclassical limit of  $\hat{H}$  established
- Final picture equivalent to background independent, Hamiltonian “floating” lattice gauge theory

# Open Questions

- Proposal for removing graph dependence (preservation), non separability, controlling fluctuations of all dof:  
Algebraic Quantum Gravity (AQG) [Giesel, T.T. 06]
- Implementation of classical  $\mathcal{N}$  symmetry (Wilson renormalisation)
- $\hat{H}$  stable coherent states?
- Better understanding of validity/physics of dust, other types of matter?
- Scrutinise LQG/AQG by further consistency checks

# Open Questions

- Proposal for removing graph dependence (preservation), non separability, controlling fluctuations of all dof:  
Algebraic Quantum Gravity (AQG) [Giesel, T.T. 06]
- Implementation of classical  $\mathcal{N}$  symmetry (Wilson renormalisation)
- $\hat{H}$  stable coherent states?
- Better understanding of validity/physics of dust, other types of matter?
- Scrutinise LQG/AQG by further consistency checks

# Open Questions

- Proposal for removing graph dependence (preservation), non separability, controlling fluctuations of all dof:  
Algebraic Quantum Gravity (AQG) [Giesel, T.T. 06]
- Implementation of classical  $\mathcal{N}$  symmetry (Wilson renormalisation)
- $\hat{H}$  stable coherent states?
- Better understanding of validity/physics of dust, other types of matter?
- Scrutinise LQG/AQG by further consistency checks

# Open Questions

- Proposal for removing graph dependence (preservation), non separability, controlling fluctuations of all dof:  
Algebraic Quantum Gravity (AQG) [Giesel, T.T. 06]
- Implementation of classical  $\mathcal{N}$  symmetry (Wilson renormalisation)
- $\hat{H}$  stable coherent states?
- Better understanding of validity/physics of dust, other types of matter?
- Scrutinise LQG/AQG by further consistency checks

# Open Questions

- Proposal for removing graph dependence (preservation), non separability, controlling fluctuations of all dof:  
Algebraic Quantum Gravity (AQG) [Giesel, T.T. 06]
- Implementation of classical  $\mathcal{N}$  symmetry (Wilson renormalisation)
- $\hat{H}$  stable coherent states?
- Better understanding of validity/physics of dust, other types of matter?
- Scrutinise LQG/AQG by further consistency checks

# Outlook

If LQG/AQG pass consistency tests then:

- All LQG techniques developed so far can be imported to phys. HS level!
- Physical semiclassical techniques to make contact with standard model
- phys. Hamiltonian defines  $S$  – Matrix, scattering theory, Feynman rules
- conservative system, hence possible improvement of vacuum problem in QFT on time dep. backgrounds (cosmology)

# Outlook

If LQG/AQG pass consistency tests then:

- All LQG techniques developed so far can be imported to phys. HS level!
- Physical semiclassical techniques to make contact with standard model
- phys. Hamiltonian defines  $S$  – Matrix, scattering theory, Feynman rules
- conservative system, hence possible improvement of vacuum problem in QFT on time dep. backgrounds (cosmology)

# Outlook

If LQG/AQG pass consistency tests then:

- All LQG techniques developed so far can be imported to phys. HS level!
- Physical semiclassical techniques to make contact with standard model
- phys. Hamiltonian defines S – Matrix, scattering theory, Feynman rules
- conservative system, hence possible improvement of vacuum problem in QFT on time dep. backgrounds (cosmology)

# Outlook

If LQG/AQG pass consistency tests then:

- All LQG techniques developed so far can be imported to phys. HS level!
- Physical semiclassical techniques to make contact with standard model
- phys. Hamiltonian defines  $S$  – Matrix, scattering theory, Feynman rules
- conservative system, hence possible improvement of vacuum problem in QFT on time dep. backgrounds (cosmology)

**HAPPY BIRTHDAY, ABHAY !!!**